A model of information filtration by comparison of randomly chosen sources

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Abstract

We study a simple model of the stochastic information filtering, in a randomly organized information system. For simplest versions of the model it appears to be possible to describe the filtering dynamics in terms of the master equations. Exact analytical results for these equations and results of numerical investigation of the dynamical features of the filter are presented.

1 Introduction

The main task of many modern information technologies is to increase the velocity of information processing. The important theoretical problem in this region is to find the basic principles of efficient organization of information flows. One of the possible approaches for its solution is developed in the framework of mathematical modeling of universal self-organization mechanisms in complex dynamical systems.

Often in this investigation, it appears to be useful to have some physical picture of phenomena under consideration. Usually, one can imagine the information system, as a net of the connected information sources. Its physical role is to realize interaction between suppliers and users of information. Although this interaction forms laws of the system evolution, in many situations it can be presented effectively

as some specific of elementary dynamical rules for the sources in the model of information net.

Recently a model of this type was proposed by A. Capocci, F.Slanina and Y.-C. Zhnag [1]. It describes ranking and filtration of information, being widespread procedures of information processing. The investigated in [1] model is a 1-dimensional stochastic dynamical system with nearest neighbor interaction.

The structure of many information system in reality is very complex, and frequently, it can be better presented by a random graph than by 1-dimensional chain. Experience in studies of physical phenomena shows that normally the interaction structure influences essentially on the system behavior. Therefore it seems to be interesting to investigate the random neighbor version of the proposed in [1] model ¹. We present in this paper the numerical and analytical results obtained for information filter dynamics of such a kind.

2 Formulation of model

We will consider a version of the filter dynamics without a fixed interaction topology. The simplest modification of the model [1] could be formulated in the following way. There are n elements with characteristics called qualities and being a number from interval [0,1]. The initial state of the system is chosen at random. The state at the time point t+1 is obtained as follows. One chooses at time point t some 2 random elements and at the time point t+1 the qualities of both elements will be equal to the same number Q(t+1). We consider two versions of the model - model A (MA) and model B (MB). In the MA the elements can be chosen arbitrarily. In the MB the quality of the chosen elements must be different, and when there are not different elements its state does not change. If the qualities of elements chosen at time t are q and p then Q(t+1)=q with probability $\frac{q}{p+q}=\mu$ and Q(t+1)=p with probability $\frac{p}{p+q}=\lambda$. If the qualities of the chosen elements were the same, it does not change. Thus, in both models the state of the system becomes stable if all of its elements are of the same quality. The dynamics of proposed models can be investigated by exact analytical methods. We demonstrate it for the initial condition of a special case. We suppose that in the initial state lelements

¹The statement of this problem was formed in discussion with Y.-C. Znahg of the presented in [1] principles of information filtering.

are of the quality p and m = n - l of the quality q. This dynamical situation arouses in the general case until the system reach the stable state. The qualitative description of dynamics of such kind can be obtained with the help of master equations.

3 Numerical results

We have chosen the MA model for numerical analysis as it is more complicated for analytical studies. Numerical experiment have been performed with the help of a distributed C program running on Solaris cluster, calculating all possible combinations of model variations, initial distributions, and number of elements; with following analysis of obtained results in Mathematica to detect model behaviour patterns.

The main characteristic of MA dynamics appears to be the scale invariance in respect to the number of elements N, which is independent from the model variations and initial conditions. We studied the system's average value of elements quality

$$A(t,N) \equiv \langle Q(t,N) \rangle$$

where Q(t, N) denotes the mean value of the system's element quality at the time point t and < ... > means averaging over ensembles of system evolution. Our results show that we can write A(t, N) in the following form:

$$A(t,N) = A_{inv} \left(\frac{t}{N}\right) \tag{1}$$

if the average values A(0,N) given by initial conditions are independent from N. Dynamics of A(t,N) for the MA model with two types of initial states are shown on the fig. 1. We considered the systems with initial states equidistributed on the interval [0,1] (A(0,N)=0,5) and initial states with 10% elements of the quality 0,9 and 90% of elements of the quality 0,1 $(A(0,N)=0,1\cdot0,9+0,9\cdot01=0,18)$. For large systems with N>100 we obtained a close correspondence between numerical results for A(t,N) and A(t,N) and A(t,N) and A(t,N) are quidistributed initial condition the correspondence is shown on the fig. 2. For this type of initial condition we have observed

$$A_{inv}\left(\frac{t}{N}\right) = 1 - \left(2 + \frac{t}{N}\right)^{-1}. (2)$$

For small N there is a certain deviation of experimental curve A(t, N) from (1), (see fig. 3), which is caused by strong effect of random fluctuations of systems with small number of elements. We have also calculated the average deviation from the average value

$$D(t,N) = \sqrt{\langle (Q(t,N) - A(t,N))^2 \rangle}.$$

The typical curve for D(t, N) is presented on the fig. 4.

For initial state equidistributed on interval [0,1] we investigated the averaged value of number $\mathcal{N}(t,N)$ of different elements in the system The obtained curve for $\mathcal{N}(t,N)/N$ seems to be very well approximated by the function S(t,N) = N/(t+N) (see fig. 5). We modified dynamics of the MA by adding Darwin selection to the system. By using the model of natural selection mechanism proposed in [2], it was presented as follows: at each time step the state of the system is changed by the MA rules and after that the element with lowest quality in the system is replaced by the one with an arbitrarily chosen quality in the interval [0,1].

The influence of the selection mechanism on the system dynamics can be analyzed by comparison of the process in usual MA and MA with Darwin selection for the same initial conditions as it is presented on the fig. 6 for a system with elements having one of two quality values and different initial distributions of "bad" (quality 0.9) and "good" (quality 0.1) elements. We see that the Darwin selection speeds up the filtration process for A(0,N) < 0,5 and slows it down for A(0,N) > 0,5.

4 Master equations

We denote the state of the system having k elements of the quality p and n-k elements of the quality q at the time point t as $\{k;t\}$. Let $P_k(t)$ be the probability of the state $\{k;t\}$. If at the time point t+1 the system state is $\{k,t+1\}$ then it was in one of the states $\{k-1;t\}$, $\{k;t\}$ or $\{k+1;t\}$ at the previous moment t. The probability $P_{2,0}(k)$ to choose two elements of quality p in the state $\{k;t\}$ is

$$P_{2,0}(k) = \frac{k(k-1)}{n(n-1)}.$$

The probability $P_{0,2}(k)$ to choose two elements of quality q is

$$P_{0,2}(k) = \frac{(n-k)(n-k-1)}{n(n-1)}.$$

The probability $P_{1,1}(k)$ to choose one element of quality p and one of quality q is

$$P_{1,1}(k) = \frac{2k(n-k)}{n(n-1)}.$$

For $k \neq 0$, $k \neq n$, the probabilities $P(\{k;t\}, \{k';t+1\})$ of transition $\{k;t\} \longmapsto \{k';t+1\}$ in the MA can be presented as

$$P(\{k;t\},\{k;t+1\}) = P_{2,0}(k) + P_{0,2}(k),$$

$$P(\{k+1;t\},\{k;t+1\}) = P_{1,1}(k+1)\mu,$$

$$P(\{k-1;t\},\{k;t+1\}) = P_{1,1}(k-1)\lambda$$

where $\lambda = p/(p+q)$, $\mu = q/(p+q)$. For k=0 or k=n, the state $\{k;t\}$ of the system is stable, hence

$$P({0;t},{0;t+1}) = P({n;t},{n;t+1}) = 1.$$

Thus, we can write the master equation for probability $P_k(t)$ as follows:

$$P_k(t+1) = \lambda P_{k-1}(t) P_{1,1}(k-1) (1 - \delta_{k0}) +$$

$$+ \mu P_{k+1}(t) P_{1,1}(k+1) (1 - \delta_{kn}) + P_k(t) P_{0,2}(k) [(1 - \delta_{kn-1}) (1 - \delta_{kn})] +$$

$$+ P_k(t) P_{2,0}(k) [(1 - \delta_{k1}) (1 - \delta_{k0})]$$

Substituting the values of probabilities $P_{2,0}(k)$, $P_{0,2}(k)$, $P_{1,1}(k)$ we have

$$P_{k}(t+1) = \frac{2(k-1)(n-k+1)\lambda}{n(n-1)} P_{k-1}(t)(1-\delta_{k0}) + \frac{2(k+1)(n-k-1)\mu}{n(n-1)} P_{k+1}(t)(1-\delta_{kn}) + \frac{(n-k)(n-k-1)+k(k-1)}{n(n-1)} P_{k}(t).$$
(3)

By similar arguments one obtains the following master equations for the MB:

$$P_0(t+1) = \mu P_1(t) + P_0(t),$$

$$P_k(t+1) = \lambda P_{k-1}(t)(1 - \delta_{k,1}) + \mu P_{k+1}(t)(1 - \delta_{k,n-1}), \text{ for } 0 < k < n, \quad (4)$$
$$P_n(t+1) = \lambda P_{n-1}(t) + P_n(t).$$

The equations (3),(4) can be used as background of analytical investigations of MA and MB. We demonstrate how they allow one to exactly calculate the important characteristics of model's dynamics.

5 Exact results for MA

The simplest problem for MA could be to find the stationary solution of the master equation (3). In vector writing it looks like:

$$P(t) - P(t-1) = -\frac{2}{n(n-1)}AVP(t-1) \quad \text{for } t > 0$$
 (5)

where P(t) is the vector with components $\{P(t)\}_k = P_k(t), k = 0, 1, ..., n$, and A, V are the matrices:

$$V_{ij} = \delta_{ij}i(n-i), \quad A_{ij} = \delta_{ij} - \lambda \delta_{ij+1} - \mu \delta_{ij-1}$$

The stationary solution P(t) = P of (5) satisfies the equation:

$$AVP = 0 (6)$$

Since det $A \neq 0$, it follows from (6) that VP = 0. Hence, the solution of (6) can be written as:

$$\{P\}_i = \delta_{i0}\rho + \delta_{in}(1-\rho), \quad 0 \le \rho \le 1 \tag{7}$$

Now, the problem is to find ρ as the function of initial probability distribution $P_k(0) \equiv \{P_0\}_k$. If we denote

$$(I - \frac{2}{n(n-1)}AV) = G$$

then in virtue of (5), the vector P is expressed through the vector P_0 in the following way

$$P = G_{as}P_0$$
, where $G_{as} = \lim_{t \to \infty} G^t$. (8)

It follows from (7),(8) that the matrix G_{as} must be of the form

$$\{G_{as}\}_{ik} = \delta_{i0}x_k + \delta_{in}y_k$$

The matrix G has the property:

$$\sum_{i=0}^{n} \{G\}_{ik} = 1 \quad \text{for} \quad 0 \le k \le n,$$

and the same equality must be fulfilled for G_{as} too, hence $x_k + y_k = 1$. Therefore

$$\{G_{as}\}_{ik} = \delta_{i0}x_k + \delta_{in}(1 - x_k).$$

The vectors $P^{(0)}$, $P^{(n)}$ with components $P_k^{(0)} = \delta_{0k}$, $P_k^{(n)} = \delta_{nk}$ are the eigen ones for matrix G: $GP^{(0)} = P^{(0)}$, $GP^{(n)} = P^{(n)}$. Hence, $G_{as}P^{(0)} = P^{(0)}$, $G_{as}P^{(n)} = P^{(n)}$, and

$$x_0 = 1, \quad x_n = 0.$$
 (9)

Taking into account that

$$G_{as}G = G_{as}$$

we obtain:

$$x_l G_{lk} = x_k.$$

Thus, we have to solve the equation xAV = 0 which can be written in components as

$$x_j - \lambda x_{j+1} - \mu x_{j-1} = 0, \quad 0 < j < n$$
 (10)

The equations (10) with boundary condition (9) coincide with ones of the classical problem of the player's losing 2 (see for example eq. (2.1),(2.2) of capital XIV in [3]). The solution of (9), (10) has the form

$$x_k = \mu^k \frac{\lambda^{n-k} - \mu^{n-k}}{\lambda^n - \mu^n} = \frac{\omega^k - \omega^n}{1 - \omega^n}.$$

Here, we used the convenient notation $\omega = \mu/\lambda$.

In virtue of (7),(8), the parameter α defining the stationary solution of the master equation can be expressed in terms of initial probability distribution $P_k(0)$ as follows

$$\rho = \sum_{k=0}^{n} P_k(0) x_k = \frac{P(\omega) - \omega^n}{1 - \omega^n}$$
(11)

We denoted $P(\omega)$ the generating function

$$P(\omega) = \sum_{k=0}^{n} P_k(0)\omega^k.$$

For the homogeneous initial distribution $P_k(0) = \frac{1}{n+1}$ we have:

$$P(\omega) = \frac{1}{n+1} \frac{1 - \omega^{n+1}}{1 - \omega}, \quad \rho = \frac{1 - (1+n)\omega^n + n\omega^{n+1}}{(n+1)(1-\omega)(1-\omega^n)}.$$

²It was pointed out us by A.Capocci.

For large n we get

$$\rho = 1 - \frac{\omega}{(\omega - 1)(n + 1)} + \omega^{-n} n Q_1(n), \text{ if } \omega > 1,$$

$$\rho = \frac{1}{(n + 1)(1 - \omega)} + \omega^n Q_2(n), \text{ if } \omega < 1,$$

$$\rho = \frac{1 + e^x(x - 1)}{x(e^x - 1)} + \frac{Q_3(n)}{n}, \text{ if } \omega = 1 + \frac{x}{n}.$$

Here, the function $Q_1(n)$, $Q_2(n)$, $Q_3(n)$ have finite limits for $n \to \infty$.

6 Exact solutions for MB

Now, we consider the master equations for MB. For the system with elements of qualities p and q this model can be considered as a reformulated classical player's losing model [3]. Let us denote P(z, u) the generating function of probability distribution $P_k(t)$:

$$P(z, u) = \sum_{k=0}^{n} \sum_{t=0}^{\infty} P_k(t) z^k u^t,$$

and will use the notations

$$A(u) = P(0, u) = \sum_{t=0}^{\infty} P_0(t)t^u, \quad B(u) = \frac{1}{n!} \frac{\partial^n}{\partial z^n} P(z, u) \bigg|_{z=0} = \sum_{t=0}^{\infty} P_n(t)t^u.$$
(12)

Then the equations (4) can be rewritten for generating function as

$$\frac{P(z,u)-P(z,0)}{u} = \left(\lambda z + \frac{\mu}{z}\right)P(z,u) + \left(1-\lambda z - \frac{\mu}{z}\right)\left[A(u) + z^n B(u)\right]$$

or in an equivalent form:

$$P(z, u)[z - u(\lambda z^{2} + \mu)] = zP(z, 0) + u\left(z - \lambda z^{2} - \mu\right)[A(u) + z^{n}B(u)](13)$$

If the functions A(u), B(u) are known, the solution of equation (13) is the following

$$P(z, u) = \frac{zP(z, 0) + u(z - \lambda z^2 - \mu)[A(u) + z^n B(u)]}{z - u(\lambda z^2 + \mu)}$$

Let us denote $\alpha_1(u)$, $\alpha_2(u)$ the solutions of equation $z - u(\lambda z^2 + \mu) = 0$:

$$\alpha_1(u) \equiv \frac{1 - \sqrt{1 - 4u^2\lambda\mu}}{2u\lambda}, \quad \alpha_2(u) \equiv \frac{1 + \sqrt{1 - 4u^2\lambda\mu}}{2u\lambda},$$

then substituting $z = \alpha_1 = \alpha_1(u)$ and $z = \alpha_2 = \alpha_2(u)$ in (13) we obtain two equations for A = A(u), B = B(u)

$$A + \alpha_1^n B = \frac{\alpha_1 P_1}{u(\lambda \alpha_1^2 - \alpha_1 + \mu)}, \quad A + \alpha_2^n B = \frac{\alpha_2 P_2}{u(\lambda \alpha_2^2 - \alpha_2 + \mu)}$$
 (14)

Here we denoted $P_i = P(\alpha_i(u), 0)$, i = 1, 2. The solution of the equations (14) has the form:

$$A(u) = \frac{\alpha_2^n P_1 - \alpha_1^n P_2}{(1 - u)(\alpha_2^n - \alpha_1^n)}, \quad B(u) = \frac{P_1 - P_2}{(1 - u)(\alpha_1^n - \alpha_2^n)}$$
(15)

Thus, we have obtained

$$P(z,u) = \frac{zP(z,0)}{[z - u(\lambda z^2 + \mu)]} +$$

$$+ \frac{u(z - \lambda z^2 - \mu)[(z^n - \alpha_2^n)P_1 - (z^n - \alpha_1^n)P_2]}{(1 - u)[z - u(\lambda z^2 + \mu)](\alpha_1^n - \alpha_2^n)}$$
(16)

It is the presentation in terms of generating function of exact solution for the master equation (4) for the finite system with n elements.

For the limit of infinite system $n \to \infty$ we obtain from (13) more simple equation:

$$P(z,u)[z - u(\lambda z^{2} + \mu)] = zP(z,0) + u(z - \lambda z^{2} - \mu)A(u).$$
 (17)

It follows from (17) that

$$A(u) = \frac{P(\alpha(u), 0)}{1 - u},$$

where $\alpha(u) = \alpha_1(u)$ is the analytical at the point u = 0 solution of equation $z - u(\lambda z^2 + \mu) = 0$. Thus, the solution of (17) has the form:

$$P(z,u) = \frac{z(1-u)P(z,0) + u(z - \lambda z^2 - \mu)P(\alpha(u),0)}{(1-u)[z - u(\lambda z^2 + \mu)]}$$
(18)

The moments $M^{(n)}(t)$ of the considered distribution function $P_n(t)$ defined as

$$M^{(s)}(t) = \sum_{k=0}^{n} P_k(t)k^s,$$

can be found by differentiating the generating function P(z, u). Particularly,

$$M^{(1)}(t) = \frac{1}{t!} \frac{\partial}{\partial z} \frac{\partial^t}{\partial u^t} P(z, u) \bigg|_{z=1, u=0},$$

$$M^{(2)}(t) = \frac{1}{t!} \frac{\partial^2}{\partial z^2} \frac{\partial^t}{\partial u^t} P(z, u) \bigg|_{z=1, u=0} + M^1(t)$$

Hence for $M^1(t)$, $M^2(t)$ can be expressed in terms of power series of coefficients of the functions

$$m_1(u) = \frac{\partial}{\partial z} P(z, u) \bigg|_{z=1}, \quad m_2(u) = \frac{\partial^2}{\partial z^2} P(z, u) \bigg|_{z=1}$$

Since

$$P(1,0) = 1, \quad \frac{\partial}{\partial z} P(z,0) \bigg|_{z=1} = M^{(1)}(0),$$
$$\frac{\partial^2}{\partial z^2} P(z,0) \bigg|_{z=1} = M^{(2)}(0) + M^{(1)}(0).$$

It follows from (18) that for infinite system

$$m_1(u) = \frac{M^{(1)}(0)}{1-u} + \frac{(2\lambda - 1)u(1 - P(\alpha, 0))}{(1-u)^2},$$

$$m_2(u) = \frac{M^{(2)}(0)}{1-u} + \frac{2(2\lambda - 1)uM^{(1)}(0)}{(1-u)^2} + \frac{2u(1 + 4\lambda^2 u - \lambda(1+3u))(1 - P(\alpha, 0))}{(1-u)^3},$$

and

$$M^{(1)}(t) = M^{(1)}(0) + \frac{1-\lambda}{\lambda} \chi_1 + (2\lambda - 1) \left[1 - \chi_0\right] t + \frac{\left[4\lambda(1-\lambda)\right]^{\frac{t}{2}}}{t^{\frac{3}{2}}} \mathcal{O}_1(t)),$$

$$M^{(2)}(t) = M^{(2)}(0) - \frac{1-\lambda}{\lambda} \chi_1 - \frac{(1-\lambda)^2}{\lambda^2} \chi_2 +$$

$$+ \left[2(2\lambda - 1) \left(M^{(1)}(0) + \frac{(1-\lambda)}{\lambda} \chi_1\right) - 4\lambda^2 (1 - \chi_0)\right] t +$$

$$+ (1-2\lambda)^2 (1-\chi_0) t^2 + \frac{\left[4\lambda(1-\lambda)\right]^{\frac{t}{2}}}{t^{\frac{3}{2}}} \mathcal{O}_2(t)$$

Here $\mathcal{O}_i(t) = 0$, i = 1, 2 are limited in the region of large t functions:

$$|\mathcal{O}_i(t)| < \mathcal{C}$$
 for $t >> 1$,

where C is a constant, and

$$\chi_0 = P(\omega, 0), \quad \chi_1 = \frac{d}{dx} P(x, 0)|_{x=\omega}, \quad \chi_2 = \frac{d^2}{dx^2} P(x, 0)|_{x=\omega}.$$

We call filtering to be finished at time point t if at this moment the quality of all the element became first to be equal. Let us denote $P^{(p)}(t)$, $P^{(q)}(t)$ the probabilities that filtering is finished at time point t with quality of all the element being equal to p and q consequently. We have

$$P_0(t+1) = P^{(p)}(t+1) + P_0(t), P_n(t+1) = P^{(q)}(t+1) + P_n(t).$$

These relations can be rewritten on terms of generating functions as

$$\mathcal{P}^{(p)}(u) = A(u)(1-u), \quad \mathcal{P}^{(q)}(u) = B(u)(1-u).$$

where A(u), B(u) are defined in (12) and $\mathcal{P}^{(p)}(u)$ ($\mathcal{P}^{(p)}(u)$) is the generating function of probabilities $P^{(p)}(t)$ ($P^{(p)}(t)$):

$$\mathcal{P}^{(p)}(u) = \sum_{t=0}^{\infty} P^{(p)}(t)u^t, \quad \mathcal{P}^{(q)}(u) = \sum_{t=0}^{\infty} P^{(q)}(t)u^t.$$

In virtue of (15), it follows that

$$\mathcal{P}^{(p)}(u) = \frac{\alpha_1^n P_2 - \alpha_2^n P_1}{\alpha_1^n - \alpha_2^n}, \quad \mathcal{P}^{(q)}(u) = \frac{P_1 - P_2}{\alpha_1^n - \alpha_2^n}$$
(19)

Hence, for generating function $\mathcal{P}(u)$ of probability $P(t) = P^{(p)}(t) + P^{(q)}(t)$ that filtering is finished at the moment t (called in [1] search time) we obtain:

$$\mathcal{P}(u) \equiv \sum_{t=0}^{\infty} P(t)u^{t} = \frac{(1 - \alpha_{2}^{n})P_{1} - (1 - \alpha_{1}^{n})P_{2}}{\alpha_{1}^{n} - \alpha_{2}^{n}}$$
(20)

For the mean value T of the filtering (search) time we obtain:

$$T = \frac{\partial}{\partial u} \mathcal{P}(u)|_{u=1} = \frac{M^{(1)}(0)}{\mu - \lambda} - n \frac{1 - P(\omega)}{(\mu - \lambda)(1 - \omega^n)}$$
(21)

This result agrees with known formula for mean time of play in the problem of player's losing (see for example capital XIV, formula (3.5)

in [3]). For the homogeneous initial probability distribution $P_k(0) = 1/(n+1)$ we have:

$$T = n \frac{(1+\omega)(1-\omega^n) - n(1-\omega)(1+\omega^n)}{2(n+1)(\mu-\lambda)(1-\omega)(1-\omega^n)},$$
 (22)

and for large n: $T = n/2|\mu - \nu| + \mathcal{O}(n), |\mathcal{O}(n)| < C$, where C is a constant.

The stationary solution of the master equation for the MB can be found as the residue of the pole in the point u = 1 of P(z, u). For the generating function $P_{st}(z) = \sum_{k=0}^{n} P_k z^k$ of stationary distribution P_k , we obtain from (16) the following result

$$P_{st}(z) = \text{res}_{u=1}[P(z, u)] = \frac{(z^n - 1)P(\omega) - (z^n - \omega^n)}{\omega^n - 1}.$$
 (23)

Comparing (7),(11) and (23), we see that the stationary solutions of the master equations for MA and MB coincide.

7 Conclusion

We studied simple processes of information filtering generated by consequent comparisons of two randomly chosen elementary information units. For the simplest version of MB the mean search time T is given by (22). It follows directly from dynamical rules that the search time for MA must be larger. The search time is dependent on initial distribution too. Nevertheless, the obtained numerical and analytical results shows that for large system the search time T and the number N of the elements obey the relation T/N = C, where C is independent on T and N and depend on initial distribution and qualities of elements only. For the MB with elements of quality p, q and homogeneous initial distribution $P_k(0) = 1/(n+1)$, quantity C looks like C = (p+q)/2|p-q| and and becomes large for small |p-q|.

The introduced in [1] characteristic of information filtering called efficiency R is defined as the rank of selected value in the starting configuration. The efficiency for the MA and MB with elements of the quality p and q can be estimated by using the stationary solution of the master equations (7). If p > q the quantity $\mathcal{R} = (1 - \rho)/\rho$ is the ratio of probabilities to select the elements with qualities p and q. It can be considered as a measure of filtration efficiency. For homogeneous initial distribution and large number N of the system elements $\mathcal{R} = N(1 - \omega)$.

The information filter investigated in [1] is characterized for large systems by search time $T \sim N^2$ and efficiency $R \sim \ln N$. Comparison these results with our ones shows that the information filtration respecting one dimensional organization of information units appears to be slower and less effective as one based on the random choosing algorithm.

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